





Three-dimensional Shape Explains Star Formation Mystery of California and Orion A

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Abstract

The new Gaia data release (EDR3) with improved astrometry has opened a new era in studying our Milky Way in fine detail. We use Gaia EDR3 astrometry together with 2MASS and WISE photometry to study two of the most massive molecular clouds in the solar vicinity: Orion A and California. Despite having remarkable similarities in the plane of the sky in terms of shape, size, and extinction, California has an order of magnitude lower star formation efficiency. We use our state-of-the-art dust mapping technique to derive the detailed three-dimensional (3D) structure of the two clouds, taking into account both distance and extinction uncertainties, and a full 3D spatial correlation between neighboring points. We discover that, despite the apparent filamentary structure in the plane of the sky, California is a flat 120 pc-long sheet extending from 410 to 530 pc. We show that not only Orion A and California differ substantially in their 3D shapes, but also Orion A has considerably higher density substructures in 3D than California. This result presents a compelling reason why the two clouds have different star formation activities. We also demonstrate how the viewing angle of California can substantially change the cloud's position in the Kennicutt–Schmidt relation. This underlines the importance of 3D information in interpreting star formation relations and challenges studies that rely solely on the column density thresholds to determine star formation activities in molecular clouds. Finally, we provide accurate distance estimates to multiple lines of sight toward various parts of the two clouds.

Unified Astronomy Thesaurus concepts: [Interstellar medium \(847\)](#); [Interstellar dust \(836\)](#); [Molecular clouds \(1072\)](#); [Star forming regions \(1565\)](#); [Star formation \(1569\)](#)

1. Introduction

Plane-of-the-sky studies have provided valuable information about the star-forming interstellar medium (ISM). However, our knowledge of the true, three-dimensional (3D) structure of molecular clouds remains limited, which hampers our understanding of how the clouds evolve and form stars. But now, thanks to the unprecedented astrometry from the Gaia mission (Gaia Collaboration et al. 2016), studying the 3D structure of molecular clouds has become possible (e.g., Großschedl et al. 2018; Leike et al. 2020; Rezaei et al. 2020; Zucker et al. 2020; Kainulainen et al. 2022). In this Letter, we compare the 3D shapes of two molecular clouds and demonstrate how that knowledge has a fundamental effect in understanding their star formation activities.

The California and Orion A molecular clouds are two of the most massive giant molecular clouds (GMCs) within 500 pc from the Sun. The two clouds are located relatively near one another and show comparable kinematics and similar filamentary shapes and sizes in the plane of the sky (Lada et al. 2009). However, their star formation activities differ significantly, with Orion A having an order of magnitude higher star formation rate than California (Lada et al. 2009). The difference triggers an immediate question about what regulates star formation and what role the individual cloud's properties like shape and size play in setting its star formation activity.

To deepen our understanding of the contrasting star formation rates in California and Orion A, we investigate their 3D shape by mapping the dust distributions toward the two clouds. The detailed 3D dust distribution toward Orion A was

already presented in Rezaei et al. (2020) where we revealed new information on its 3D shape, discovered a foreground dust component, and demonstrated the extending tail of the cloud to further distances as suggested by Großschedl et al. (2018). Here, we look into the 3D distribution of the dust toward California to map its 3D substructures that remain hidden in the plane-of-the-sky data.

While there are a couple of local dust maps of the Milky Way, our current work toward California and Orion A is a unique map of its kind. Leike et al. (2020) provide a 3D dust map of the solar neighborhood using Gaia DR2 and metric Gaussian variational inference that is then used by Zucker et al. (2021) to define cloud boundaries. However, Leike et al. (2020) provides data for all lines of sight (LOSs) only up to 370 pc and for some LOSs close to 500 pc, depending on the coordinates. Their results beyond 370 pc need to be treated with extreme caution, as it reaches the boundaries of the map that affects the density, shape, and distance to the clouds (Zucker et al. 2021). Apart from the coverage limits of Leike et al. (2020), our 3D dust mapping technique has the advantage of taking into account distance uncertainties; as a result, it can exploit a full data set like Gaia without the necessity to cut on noisy data. Those stars are typically located in dusty regions, which are important for mapping the dense parts of the molecular clouds. Moreover, owing to the analytical solution for our posterior calculation, approximations are minimal; therefore, our method provides accurate and reliable estimates (see Rezaei et al. 2017, 2018b, 2020 for more details).

2. Method and Data

Our 3D dust mapping technique has been extensively explained in Rezaei et al. (2017, 2018b). In addition, new



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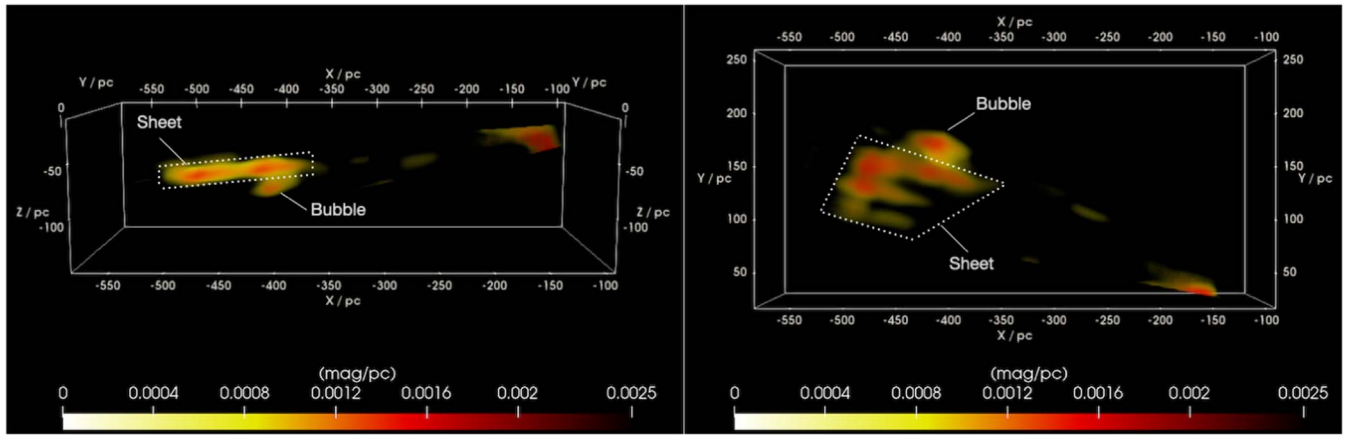


Figure 1. 3D dust density predictions for the California cloud projected from the y -axis (left panel) and the z -axis (right panel) (perpendicular to the Galactic plane). More angles can be seen in the Appendix A (Figure 6). The color represents densities in K_s -band magnitude per parsec. The Sun is at $(0,0,0)$ and X increases toward the Galactic Cent. The dotted lines demonstrate what we refer to as a sheet. The predictions are made on regular grids for every 0.5° in the Galactic l and b , and every 5 pc in distance. The 3D image is then produced by applying a smoothing kernel to handle the missing pixels. In order not to produce extra smoothing from that of the method, the length scale of the smoothing kernel is chosen to be much smaller (3 pc) than the correlation length. For illustration purposes, values below 0.0005 are set to be transparent. For $N_H(\text{cm}^{-2}) = 2 \times 10^{22} A_K$, the dust density of $0.001 \text{ mag pc}^{-1}$ corresponds to a gas volume density of $\sim 6 \text{ cm}^{-3}$ (see Figure 3).

developments have been introduced in Rezaei et al. (2020). Here we briefly summarize the main aspects of our method.

Our technique consists of a nonparametric method that uses 3D positions of stars together with their LOS attenuation as the input data and predicts dust densities for arbitrary points in the same 3D space. We divide the LOS toward stars into 1D cells in order to model the integrated attenuation to each star as sum of the dust densities along its LOS by connecting all 1D cells using a Gaussian process (GP) prior. GP takes into account the neighboring correlations using a correlation length, λ . Another hyperparameter of the model is the variance, θ , which sets the amplitude of density variations. All hyperparameters are fixed based on the input data as explained in Rezaei et al. (2017, 2018b). The final resolution of the map is set by the typical separation between input stars. Apart from considering the 3D neighboring correlation, our method has the advantage of taking into account both distance and extinction uncertainties toward individual stars that result in producing robust estimates of the distance and 3D shapes of the molecular clouds. Despite common artifacts in most of the 3D dust maps, our results are devoid of discontinuity and “fingers-of-God” artifact.

The early installment of the third Gaia data release (Gaia EDR3) consists of the full astrometric solution for around 1.5 billion sources, and a significant advantage over Gaia DR2 parallaxes with 30% increase in precision (Gaia Collaboration et al. 2021). We use the 3D positions of stars from Gaia EDR3; only limiting our sample to parallax uncertainty of less than 100%. Since we are dealing with large parallax uncertainties, we use the geometric distance estimates from Bailer-Jones et al. (2021) to map the full shape of the cloud.

Similar to Rezaei et al. (2020), we use the Two Micron All-Sky Survey (2MASS; Skrutskie et al. 2006) and the Wide-Field Infrared Survey Explorer (WISE; Wright et al. 2010) photometry to calculate extinctions to individual stars using the Rayleigh–Jeans Color Excess (RJCE; Majewski et al. 2011) method, which provides us with the extinction in the K_S band. The crossmatch between 2MASS and WISE catalogs with Gaia EDR3 sources are provided on Gaia Archive.³

After calculating extinctions, we select our final sample based on the position of stars on the dereddened color–magnitude diagram in order to remove the outliers (see Rezaei et al. 2018a, 2020). For our current work toward the California cloud, $155^\circ < l < 170^\circ$ and $-14^\circ < b < -6^\circ$, we use around 160,000 stars as our final input sample. The hyperparameters of the method for the aforementioned data are cell size = 5 pc, $\lambda = 20$ pc, and $\theta = 4 \times 10^{-8} \text{ pc}^{-2}$. It is important to note that the choice of hyperparameters is not completely arbitrary. As discussed in detail in Rezaei et al. (2017, 2018b), the hyperparameters are calculated according to input data. The final resolution of the map is not set by these parameters, but rather the input data (see Appendix B for an example).

3. California in 3D

Figure 1 shows the 3D distribution of dust toward the California region from two viewing angles. Rather than being a filament perpendicular to the LOS, California is an extended flat, sheet-like structure, with a bubble at one side, where higher-density structures form within the sheet. From the face-on view, the longest length of the sheet is approximately 120 pc, with the shortest length of about 80 pc, while from the edge-on view, the width of the sheet appears very narrow (as expected from the plane-of-the-sky observations). Figure 2 demonstrates various components of the cloud at multiple distance slices: while the bubble centered at $l = 158^\circ$, $b = -10^\circ$ appears clearly at 455 pc and expands over only 25 pc in radius, some other parts of the cloud extend along the sheet for over 100 pc.

Having uncertainties on the dust density predictions is crucial for evaluating the significance and validity of different 3D substructures, as explained in Rezaei et al. (2020). Figure 3 shows dust density predictions with their uncertainties as a function of distance for three LOSs toward different parts of the cloud (also marked in Figure 2 in the plane of the sky). The red curve that is toward the lowest part of the bubble has a dominant, symmetrical overdensity at 445 pc, and there is a shallower overdensity in the foreground at 290 pc. The foreground overdensity is in agreement with an increase in reddening in the Green et al. (2019) map at the same distance,

³ <https://gea.esac.esa.int/archive/>

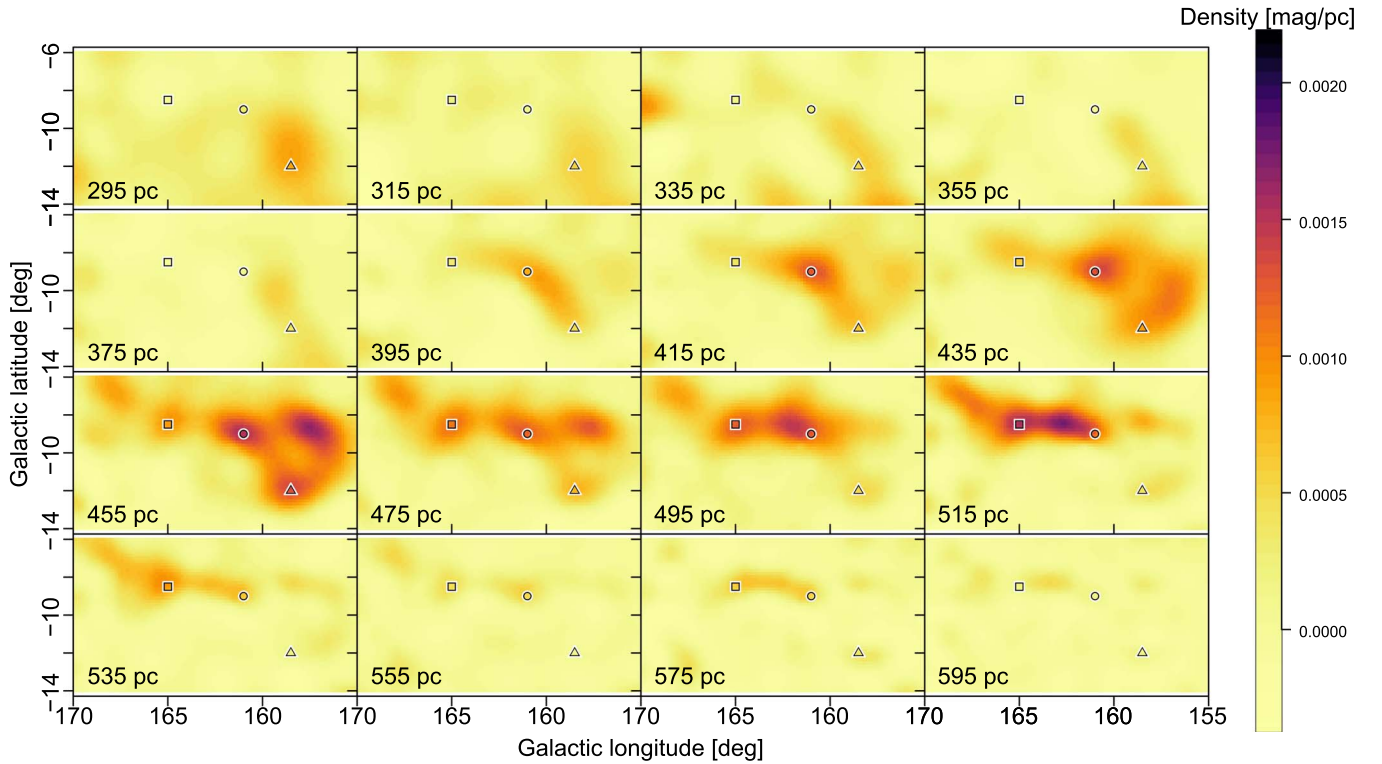


Figure 2. Dust density predictions in the plane of the sky. Each panel represents a slice through the cloud at fixed distances (every 20 pc). The bubble on the right side is apparent at the 455 pc slice, while the rest of the cloud is seen through multiple panels demonstrating the extent of the sheet, especially toward $(l, b) = (161.5, -8.5)$. The three symbols (square, circle, and triangle) represent specific LOSs along the cloud investigated in Figure 3. For illustration purposes, the image is smoothed with the scale length of 0.4° . For $N_H(\text{cm}^{-2}) = 2 \times 10^{22} A_K$, the dust density of $0.001 \text{ mag pc}^{-1}$ corresponds to a gas volume density of $\sim 6 \text{ cm}^{-3}$.

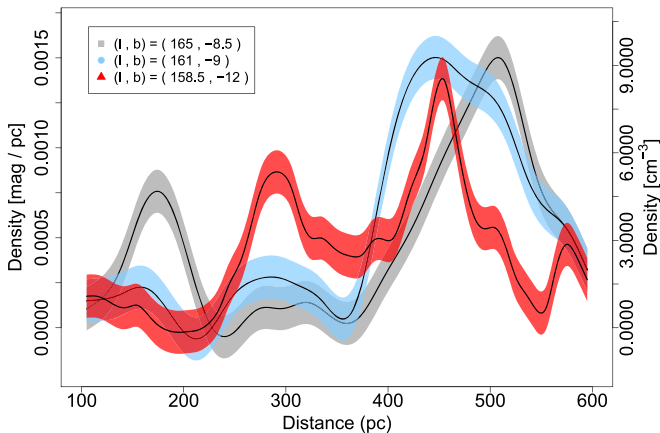


Figure 3. Dust density vs. distance for three different LOSs toward the California cloud (overplotted in Figure 2). The black line shows the mean, and the shading represents one standard deviation (also computed by the Gaussian process model). The red curve is toward the lowest part of the bubble, the blue curve represents the elongated length of the sheet, and the gray curve shows the density for an LOS toward the left (eastern) end of the cloud. The density in the secondary axis is calculated assuming $N_H(\text{cm}^{-2}) = 2 \times 10^{22} A_K$, which is an average value for the hydrogen column density in the literature (e.g., Reina & Tareghi 1973; Gorenstein 1975; Predehl & Schmitt 1995; Güver & Ozel 2009).

and an increase in density in Leike et al. (2020). The blue and gray curves, however, recover a broader distribution that represents the extent of the sheet along the LOS. While the highest peak of dust density along $(l, b) = (161, -9)$ (blue curve) is around 450 pc (as estimated by Zucker et al. 2020), it captures the most elongated part of the sheet extending from 410 to 530 pc, giving it the length of about 120 pc. There is also

a shallow overdensity at closer distances toward this LOS that is related to the Taurus cloud on the foreground. As seen from the gray curve, $(l, b) = (165, -8.5)$, the peak density of the eastern part of California is at a farther distance of 514 pc.

It is important to note that the extension of the clouds along various LOSs are driven by the input data; in fact, as explained in our previous works, our method is capable of capturing multiscale density substructures that are smaller or larger than the input correlation length. In addition, the peak densities presented here are not representative of the high-density small-scale molecular gas but an average density within the resolution of the map; therefore, much higher density cores can be located within our dust clumps. We demonstrate these points with a mock data set in Appendix B.

We can reconstruct the 2D extinction map toward California by integrating our predicted 3D dust densities along various LOSs. As can be seen from Figure 7, the final 2D projected map reaches extinction of $\sim 0.5 \text{ mag}$ in the K_s band ($\sim 5 \text{ mag}$ in A_V), and there is a clear resemblance to higher-resolution maps (e.g., CO-based maps or NICEST/NICER extinction maps).

4. California versus Orion A

Even though very similar in the plane of the sky (Lada et al. 2009), California's 3D structure appears very different from that of Orion A. Within our resolution limit, only one apparent feedback-driven substructure is seen in California (a bubble on the western side), while Orion A has several such structures. Beyond this substructure, California is distributed along a flat sheet, while Orion A is more filamentary also in 3D. The two clouds are shown together in Figure 4. The substructures in Orion A are more complex than in California, likely caused by

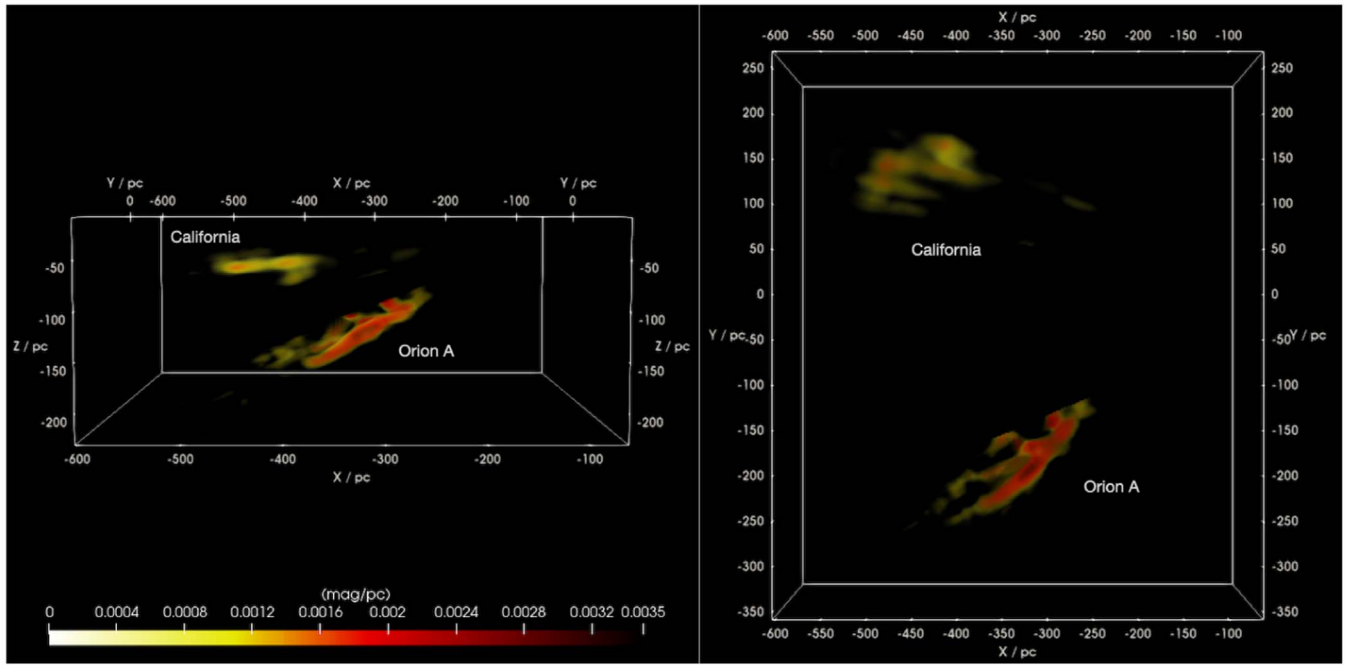


Figure 4. Two projected views of the 3D distribution of dust toward the California and Orion A clouds. The Sun is at (0,0,0), the z -axis is perpendicular to the Galactic plane, and X increases toward the Galactic center. For illustration purposes, values below 0.0005 are set to be transparent.

feedback processes and multiple episodes of star formation in the region (e.g., Schlafly et al. 2015; Rezaei et al. 2020). In addition, substructures in Orion A have higher densities than in California (see Figure 4). Since the shallow overdensities in California are spread along the flat sheet, they add up in the plane of the sky to represent extinction properties similar to those of Orion A. However, the 3D view reveals the dramatic differences between the clouds.

In addition, other studies have estimated the age of California to be around 1–3 Myr (e.g., Covey et al. 2010; Wolk et al. 2010; Imara et al. 2017), while Orion A seems to be in an older evolutionary stage with older stellar populations (5–10 Myr; e.g., Bouy et al. 2014; Zari et al. 2019) associated with its discovered foreground bubble (Rezaei et al. 2020), which seems to have triggered the next generation of star formation in the main Orion A filament in the background.

We also demonstrate that the knowledge of clouds’ 3D shape is an important factor in interpreting their star formation activities and in considering star formation relations such as the Kennicutt–Schmidt (KS) relation. Without observational biases, the mean surface density of the California cloud changes drastically depending on the viewing angle: the total mass of the cloud is the same irrespective of the viewing angle, but the surface area is an order of magnitude larger when viewed from the direction perpendicular to the Galactic plane (perpendicular to the cloud’s sheet-like morphology). The change in area is less dramatic for Orion A that is more filamentary than sheet-like. This effect has two potentially important consequences:

First, the location of California in the KS relation depends significantly on the viewing angle. The star formation rate (SFR) in a molecular cloud in the Milky Way is calculated from the number of young stellar objects (YSOs) in the region, assuming an average mass and age for the YSOs. Since the number of observed YSOs does not change with the angle, then

SFR will remain constant. As the cloud area changes, but the mass and SFR remain constant, California can move in the relation along a line with the slope of unity (see Figure 5).

How exactly the different structures of the clouds affect their location in the observed KS relation depends on the details of the observational techniques used to derive the masses and areas of the clouds. In Kainulainen et al. (2022), we showed that in particular the column density threshold used to define the clouds has a strong effect on this. In particular, using plane-of-the-sky data, Lada et al. (2013) found that there is no KS relation between the clouds in the solar neighborhood. Our result opens a door to speculate that the relation may well exist, but, depending on the size and morphology of the clouds, is not recovered from the plane-of-the-sky data. This could be either due to systematic effects in the cloud orientations, coincidence in the sampling of random orientation angles, or the details of the cloud definition process (especially thresholding; Kainulainen et al. 2022).

Second, our result also suggests a link between the 3D morphology of molecular clouds and their star formation activity. The plane-of-the-sky works indicate that the California and Orion A clouds are morphologically similar; the origin of the differing star formation activities has been attributed to differing amounts of dense gas (e.g., Kainulainen et al. 2009; Lada et al. 2013; Kainulainen & Federrath 2017) and/or environmental effects (e.g., Schlafly et al. 2015). Our work shows that the two clouds are *not* morphologically similar, which immediately raises the question about the importance of morphology. Again, providing 3D shape information for a larger sample of clouds will be vital to establish this.

Finally, we provide the most up-to-date distance estimates to various LOSs toward California and Orion A using Gaia EDR3 (Table 1). We note that the geometric distance estimates from GEDR3 parallaxes by Bailer-Jones et al. (2021) are over-estimated because their prior’s scale length is larger than our

Table 1
Distance Estimates to Orion A and California Using Gaia EDR3

Name	l ($^{\circ}$)	b ($^{\circ}$)	Distance (pc) to Peak Density	Shape/No. of Components
Orion A	206.5	-17	347 pc & 393 pc	Double peak
Orion A	209	-19	360 pc & 387 pc	Double peak
Orion A	212.5	-19	406 pc	Dominant one peak
Orion A	214.5	-20	427 pc	Dominant one peak
California	158	-10.5	437	Dominant peak, additional lower-density peak at 268 pc
California	158.5	-12	445	Dominant peak, additional lower-density peak at 290 pc
California	160.5	-9.5	455	Elongated; 410–530 pc
California	161.5	-8.5	...	Elongated; 410–530 pc, no dominant peak
California	163	-8.5	507	Dominant peak within the elongated sheet; 410–530
California	164.5	-7.5	505	Dominant one peak
California	165.3	-9	514	Elongated; 455–530 pc

Note. Our model estimates densities for fixed points in 3D space; therefore, there is no direct uncertainty assigned to our distance predictions. However, the predicted densities have full uncertainty estimates that can be propagated into distance estimates (as explained in Rezaei et al. 2020). Having done so, we estimate uncertainties for our predicted distances to Orion A and California clouds to be 10 pc.

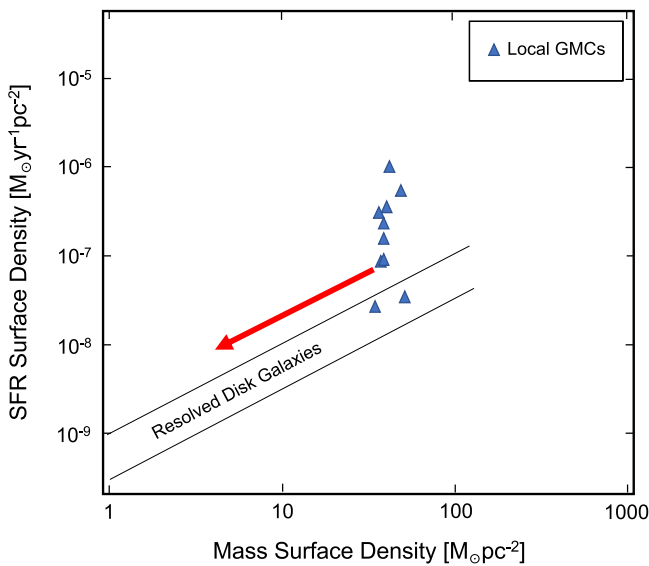


Figure 5. Star formation surface density as a function of mass surface density (Kennicutt–Schmidt relation) for the local molecular clouds from Lada et al. (2013). The red arrow shows how the location of a cloud like California would change on this diagram when seen from an angle perpendicular to the LOS.

targeted cloud distance closer than 500 pc. Therefore, after we map the full 3D shape of the cloud using all stars that match our input criteria (see Section 2), in order to get accurate distances to individual dust components along various LOSs, we only use stars from Gaia EDR3 with parallax estimates better than 20% and directly invert them to get their distances.

There is a perfect agreement between our distance estimates and distance estimates of Zucker et al. (2020) for LOSs toward where there is a prominent peak. However, where there are multiple peaks near each other along an LOS (like the foreground cloud in Orion A) or where there are no dominant peaks (e.g., elongations in California), our distance estimates differ. Although distances in Zucker et al. (2020) are predicted on higher spatial resolution than that of our work, it is possible that due to the nature of the Zucker et al. (2020) technique, multiple nearby peaks and the physical shape of the clouds remain undetected in their work. It is important to note that, as explained in Rezaei et al. (2020), our predictions are made for fixed points in 3D; therefore, there is no direct uncertainty

associated with our distances. However, our method provides the full probability distribution function for density predictions at each point that can then be propagated into distance uncertainties. Having done so, the typical uncertainty in our predicted distances is 10 pc.

5. Concluding Remarks

We have determined the 3D structure of the California molecular cloud using our advanced mapping technique that takes into account distance and extinction uncertainties to individual stars and considers correlations between neighboring points in 3D. Having the full probability density function for each predicted point, i.e., mean and standard deviation of the predicted density, allows us to verify the results; therefore, our final products are robust 3D dust maps of the molecular clouds.

We demonstrated that California is a sheet-like structure extended around 120 pc along the LOS. It consists of low-density substructures that align with the LOS resulting in higher column density in the plane of the sky, comparable to that of Orion A. The surface density of California, however, is much lower than Orion A from the viewing angle perpendicular to the Galactic plane. In addition, the substructures of the two clouds show substantial differences with Orion A being a filamentary structure with several feedback bubbles, while California has only one isolated bubble on one side and largely resembles an unperturbed sheet. The two clouds, even though similar in the plane of the sky, are remarkably different in 3D. The dramatic differences likely contribute to the different star formation efficiencies of the clouds. Our results demonstrate the importance of 3D information in understanding the star formation activities of molecular clouds.

We also showed that the 3D information plays a dramatic role in setting the mean surface density of the clouds, which can then significantly affect their appearance in scaling relations such as their location in the KS relation. This further indicates that the column density thresholds are not necessarily reliable tools to determine core star formation activities in molecular clouds. We have presented the results for California in this Letter, but plan to expand this study to more clouds in order to have a better understanding of star formation in the solar neighborhood and its connection to extragalactic studies.

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Appendix A Reconstructed Extinction

Figure 6 shows the 3D distribution of California from different viewing angles which better demonstrates the full structure of the cloud. Figure 7 shows the reconstructed extinctions (A_{K_s}) using the predicted 3D densities. For each LOS the predicted densities are summed up and multiplied by the cell sizes in order to get the projected extinctions. Despite the low resolution, the structures resemble that of the higher-resolution maps (Lada et al. 2009). Contours of $A_{K_s} = 0.3$ mag from Green et al. (2019) are shown in the figure.

Figure 8 shows the reconstructed extinctions (A_{K_s}) using the predicted 3D densities as a function of distance for three LOSs shown in Figure 3, together with those from the Green et al. (2019) map. The slope of the reddening curves represents the

amount of dust density in the corresponding distance. Therefore, an increase in the reddening in Figure 8 is equivalent to a peak in our LOS density profile in Figure 3. Overall, there is a fairly good agreement between our predicted distances to the clouds and those from Green et al. (2019); the only difference is for the red curve, where the secondary peak in our map is not recovered by Green et al. (2019).

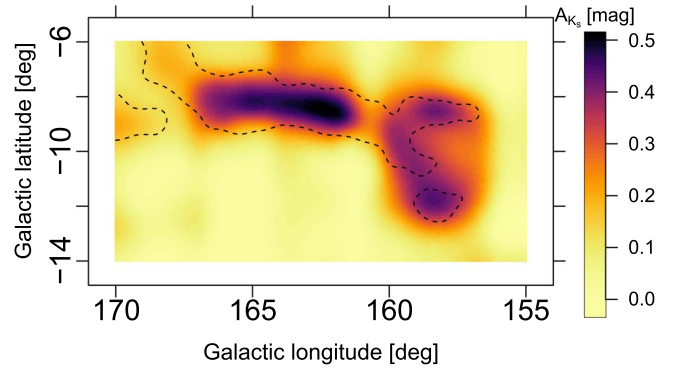


Figure 7. Reconstructed extinction using the 3D density predictions (in the K_s band). Extinctions are calculated using the sum of the predicted 3D densities along each LOS. The 2D representation of the 3D predictions nicely recover the plane-of-the-sky features of California. The dashed lines are contours of 0.3 mag extinctions from Green et al. (2019) using the “dustmaps” interface of Green (2018). For consistency, the second map is smoothed to the same resolution as our map.

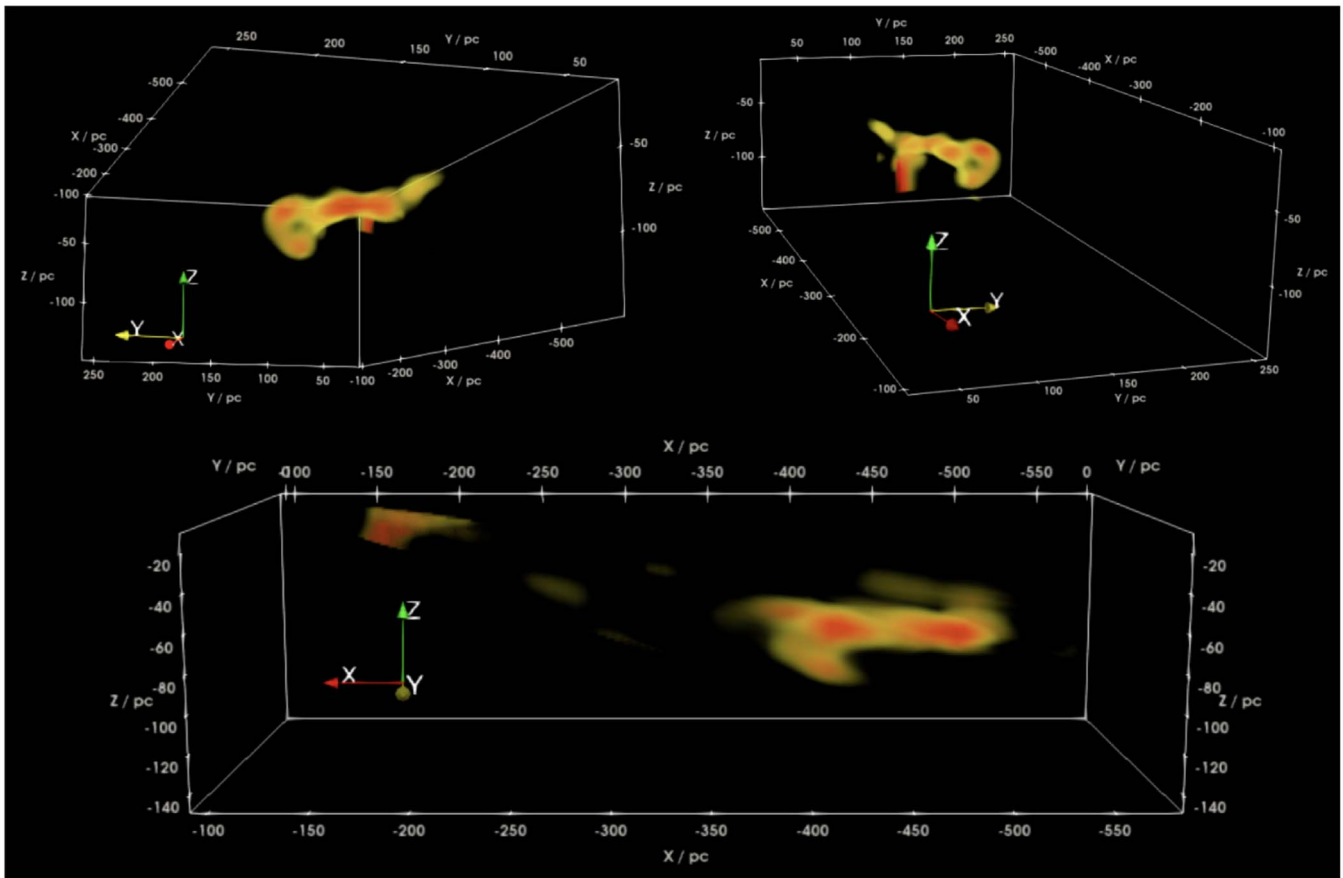


Figure 6. 3D structure of California from multiple viewing angles. Small arrows show the orientation of each angle. The bubble described previously and marked in Figure 1 is visible in multiple panels. Color scales are as in Figure 1.

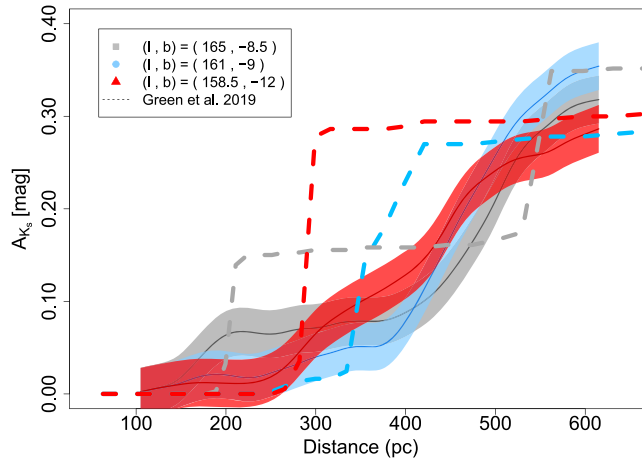


Figure 8. Extinction as a function of distance for three LOSs toward California as shown in Figure 3 (solid lines with 1σ uncertainties as shaded area), plus those from Green et al. (2019) (dashed lines). A typical uncertainty in Green et al. (2019) is about 0.02 mag in this range. The predicted distances for gray and blue LOSs agree fairly well in both maps, while the red curve from Green et al. (2019) is missing the second peak in our map.

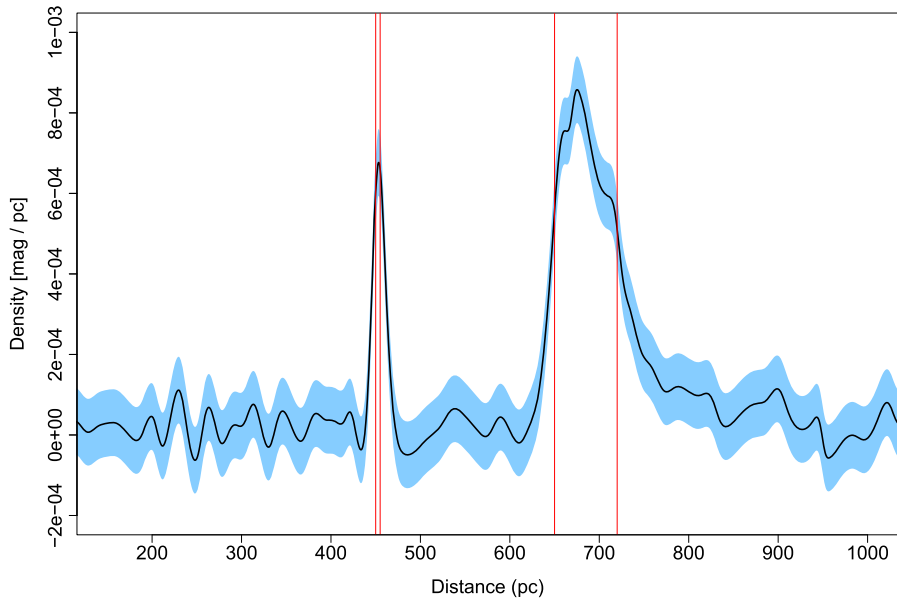


Figure 9. Density predictions as a function of distance for mock data set explained in Appendix B. The vertical red lines demonstrate the “true” distances and widths of the clouds.

Apart from the predicted distances, predicted extinction in Green et al. (2019) seems to appear in steps; allocating almost all predicted dust densities toward a cloud to a certain distance, while an increase in extinction in our map is more gradual. This could be the result of different techniques used in each work: the method used in Green et al. (2019) is not optimal for predicting multiple overdensities near one another or recovering an extended structure. Our technique, on the other hand, tends to smooth the predictions, therefore underestimating the maximum predicted density for each peak. It is important to note, however, that the amount of elongation caused by the smoothing in our method is less than the 10 pc uncertainties; thus, the predicted elongated structures in our map are real effects of the data rather than an artifact caused by the technique (see Appendix B for an example using mock data).

Appendix B Mock Data Set

Here we demonstrate the ability of the method to capture the locations and widths of the clouds quite precisely. Also, we demonstrate how the peak density predictions in the map are affected by its resolution.

We simulate mock observations using two dust clouds with different sizes and densities, and stars with similar spatial density distribution as that of Gaia. Here are the true properties of the simulated clouds:

1. Cloud 1 at 450 pc, 5 pc width, density of $2 \times 10^{-3} \text{ mag pc}^{-1}$,
2. Cloud 2 at 650 pc, 70 pc width, density of $1 \times 10^{-3} \text{ mag pc}^{-1}$.

We then calculate the LOS extinction to stars in this region and add 20% noise to both distances and extinctions in order to have a more realistic simulation. We use these as the input data and use the model to find the underlying 3D dust densities.

As can be seen in Figure 9, the model can predict the location of the clouds and their widths quite precisely. The predicted widths (elongation) of the clouds at the half-maximum are 15 pc for the first one and 85 pc for the second one, which is consistent with their true widths within ± 10 pc uncertainties. It is worth mentioning that even though the correlation length of the model was 20 pc, it recovered clouds with completely different lengths from 5 to 70 pc; indicating that the elongation of the cloud is mainly determined by the input data. Another important point about the predictions are the maximum densities: since the size of the first cloud is smaller than the scale lengths of the model, the maximum predicted density is less than half of the “true” input density. These points indicate the following: (1) Our predicted distances are quite precise. (2) The extension of the clouds is reliable and determined by the input stars (with the predicted uncertainties). (3) The maximum densities probed by the map is an average density within the resolution of the map. The black line is the mean and the blue shaded area is the standard deviation of the predictions.

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